# Dynamic modeling and performance analysis of the 2PRU-PUU parallel mechanism 

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Abstract. This paper investigates the dynamic modeling and performance analysis of the 2PRU-PUU reconfigurable parallel mechanism (RPM); here, P, R, and U denote the prismatic, revolute, and universal joints, respectively. By altering one of the rotation axes of the reconfigurable universal joint in limb 3, the mechanism can be switched into two operation modes, 1R2T and 2R1T. The authors resort to the Lagrangian equations of the first kind to derive the dynamic model of the 2PRU-PUU RPM. The optimal driving force distribution is determined to solve the problem of the non-uniqueness solution in dynamic analysis. The dynamic formulations are verified with the results obtained in ADAMS software. The dynamic manipulability ellipsoid index, which offers a quantitative assessment of the ability in manipulating the end effector, is used to assess the dynamic performance of the mechanism. Then, the distribution characteristics of the rotational and translational dynamic performance of the RPM are derived.

## 1 Introduction

Due to the positioning accuracy, excellent dynamic performance, and high stiffness (Hu et al., 2020a; Yang et al., 2022), the parallel mechanism (PM) has been widely applied in many fields, such as the Tricept (3-UPS+UP) PM (Kureková and Halaj, 2014), the Z3 head (3-PRS) (Chen et al., 2014), the serial-parallel robotic arm (3RPS+3SPR) (Hu et al., 2019), and the Exechon hybrid manipulator (Hu et al., 2020b). However, traditional PMs cannot meet the needs of multiple task requirements in the modern automation field, so PMs with multi-operation modes (Ye et al., 2014; Li and Herve, 2009), also called reconfigurable PMs (RPMs), have become a research hotspot.

The primary characteristics of RPMs are as follows: (a) fewer actuators are needed for the moving platform to realize several motion patterns, and (b) reconfiguring the PM does not require disassembly of the PM. There have been several classes and a general method for the type synthesis of RPMs (Kong, 2013; Liu and Liu, 2022), and scholars have done lots of research on reconfigurable analysis (Carbonari et
al., 2019; Kong, 2014). However, previous research only focused on the type synthesis and kinematics; research on the dynamics of RPM is quite limited.
The foundation of dynamic performance analysis for PMs is an accurate dynamic model, which is also an essential prerequisite for realizing parameter optimization and highprecision, high-efficiency control of PMs. However, dynamic analysis is more challenging than kinematic analysis because of the difficulty in describing the overall mass metrics caused by the multiple closed-loop properties of PMs. Typical dynamic modeling can be directly computed in each configuration because each configuration can be thought of as a typical PM with a fixed number and type of operation modes. Currently, the Newton-Euler method (Gan et al., 2016), the Lagrange method (Karimi Eskandary and Angeles, 2018), the principle of virtual work (Pedrammehr et al., 2018), and the Kane equation (Elgolli et al., 2019) are the main methods for the dynamic modeling of PMs.

In the reported literature, redundant actuated PMs possess many advantages, such as extended reachable workspace
(Arata et al., 2011), elimination or reduction of singularities, and increased stiffness (Hu et al., 2012; Zhao et al., 2009). However, there are also some challenges introduced by redundant actuators such as the non-uniqueness solution in dynamic analysis, which is caused by the fact that the number of equations is always less than the number of unknown variables, and thus the dynamic equation cannot be solved uniquely. In recent years, the Newton-Euler method has been used to investigate the dynamic mode of a 3-degree-of-freedom (DOF) redundantly actuated parallel manipulator by taking the flexible deformation of the limbs into account (Li et al., 2016). The forward and inverse dynamics formulations of the 3-RRR PM were developed using the Jacobian/Hessian matrices of the constraint equations, the Lagrange-D'Alembert equation, and the Hessian matrix of the kinetic energy of the manipulator (Abo-Shanab, 2020). Based on the virtual work method, an improved general dynamic formulation of the inverse and direct dynamics of the 6-UPS Gough-Stewart PM was presented (Kalani et al., 2016). Wang et al. (2019) resorted to the natural orthogonal complement based on an adaptation of screw theory to derive the dynamics model of the 2PUR-2RPU PM.

The dynamic manipulability ellipsoid (DME) index (Chen et al., 2021; Chiacchio, 2000) is used to evaluate the dynamic performance of the proposed RPM, namely the easiness of arbitrarily changing the position or orientation of the end effector of PMs. The distribution characteristics of both rotational and translational dynamic performance are obtained in an intuitive manner (Yoshikawa, 1990; Chen et al., 2017). The adopted dynamic modeling method and dynamic performance analysis will provide a basis for structural optimization and motion control of the subsequent prototype.

The main contribution of this work is the dynamic modeling of the reconfigurable PM based on the Lagrange method, which has the benefit of being able to handle the motion of the mechanism uniformly under various operation modes. The proposed dynamic model serves as a valuable tool for control algorithm and performance optimization of RPMs.

This paper is organized as follows: in Sect. 2, the structure description of the 2PRU-PUU RPM is presented. In Sect. 3, position analysis of the proposed RPM is conducted. In Sect. 4, dynamic modeling is established by the Lagrange method first, then the driving force distribution of redundant actuator PMs is proposed, and finally the correctness of the dynamic models is verified by ADAMS simulation software. In Sect. 5, the DME index is calculated within the workspace of the PM. In Sect. 6, some concluding remarks and directions for future work are summarized.

## 2 Structure description

As shown in Fig. 1, PM is composed of a moving platform, a fixed platform, two identical PRU limbs, and a PUU limb. In PRU limbs, the two-P pair has the same axis and is per-


Figure 1. The model of the 2PRU-PUU PM.
pendicular to the axis of the P pair in the PUU limb, and the direction of the R pair axis is the same as the first rotational axis of the $U$ joint connected to the moving platform. In the PUU limb, the second rotation axis of the lower $U$ joint is parallel to the first rotation axis of the upper $U$ joint, and the first rotation axis of the lower U joint is parallel to the R pair axis in the PRU limb and the second rotation axis of the upper U joint in the PUU limb. Three limbs connect the fixed and moving platform at point $P_{i}$ and $B_{i}$ with prismatic pairs and universal joints, respectively. Coordinate frames $o\{x, y, z\}$ and $o^{\prime}\{u, v, w\}$ are linked to the fixed platform and moving platform, respectively.

From the switch configuration, the RPM can evolve into two diverse configurations: one with 1 rotational DOF and 2 translational DOFs and the other with 2 rotational DOFs and 1 translational DOF (Ye et al., 2022). Additionally, since the RPM has 4 DOFs in the switch configuration, it requires a minimum of four actuators to operation completely. Therefore, the RPM will be redundantly actuated in both operation modes.

## 3 Position analysis

Kinematic and dynamic analyses are built on the inverse position solution of PMs. For the convenience of analysis, the kinematics model of 2PRU-PUU RPM shown in Fig. 2 is drawn according to the mechanism characteristics. The coordinate vector of point $P_{i}(i=1,2,3)$ in the fixed coordinate frame can be expressed as

$$
\begin{align*}
& \boldsymbol{r}_{P_{1}}=\left[\begin{array}{lll}
-d_{1} & 0 & 0
\end{array}\right]^{\mathrm{T}} \\
& \boldsymbol{r}_{P_{2}}=\left[\begin{array}{lll}
d_{2} & 0 & 0
\end{array}\right]^{\mathrm{T}}  \tag{1}\\
& \boldsymbol{r}_{P_{3}}=\left[\begin{array}{lll}
0 & d_{3} & 0
\end{array}\right]^{\mathrm{T}}
\end{align*}
$$

In the moving coordinate frame, the coordinate vector of point $B_{i}(i=1,2,3)$ can be written as

$$
\left\{\begin{align*}
\boldsymbol{r}_{B_{1}^{\prime}} & =[-a \tag{2}
\end{align*} 0000\right]^{\mathrm{T}} .
$$

The coordinate vector of point $B_{i}(i=1,2,3)$ in the fixed coordinate frame can be written as
$\boldsymbol{r}_{B_{i}}=\mathbf{R} \boldsymbol{r}_{B_{i}^{\prime}}+\boldsymbol{r}_{p i}$,
where $\mathbf{R}$ denotes the rotation matrix from the moving coordinate frame to the fixed coordinate frame.
$\mathbf{R}=\mathbf{R}(y, \theta) \mathbf{R}(x, \varphi)=\left[\begin{array}{ccc}c \theta & s \varphi s \theta & s \theta c \varphi \\ 0 & c \varphi & -s \varphi \\ -s \varphi & c \theta s \varphi & c \theta c \varphi\end{array}\right]$
Substituting Eqs. (1) and (2) into Eq. (3), $\boldsymbol{r}_{B_{i}}$ can be obtained as follows:

$$
\left\{\begin{array}{l}
\boldsymbol{r}_{B_{1}}=\left[\begin{array}{lll}
x-a c \theta & 0 & z+a s \theta
\end{array}\right]^{\mathrm{T}}  \tag{5}\\
\boldsymbol{r}_{B_{2}}=\left[\begin{array}{lll}
x+a c \theta & 0 & z-a s \theta
\end{array}\right]^{\mathrm{T}} \\
\boldsymbol{r}_{B_{3}}=\left[\begin{array}{lll}
x+a s \theta s \varphi & a c \varphi & z+a c \theta s \varphi
\end{array}\right]^{\mathrm{T}}
\end{array}\right.
$$

Using the constraint equation $L_{i}=\left|\boldsymbol{r}_{B_{i}}-\boldsymbol{r}_{p i}\right|$, three equations can be obtained as follows:

$$
\left\{\begin{array}{l}
d_{1}=a c \theta-x \pm \sqrt{L_{1}^{2}-(z+a s \theta)^{2}}  \tag{6}\\
d_{2}=a c \theta+x \pm \sqrt{L_{2}^{2}-(-z+a s \theta)^{2}} \\
d_{3}=a c \varphi \pm \sqrt{L_{3}^{2}-(a s \theta s \varphi+x)^{2}-(z+a c \theta s \theta)^{2}}
\end{array}\right.
$$

where the " $\pm$ " notation represents two distinct assembly modes for the limbs, which are selected as " + " according to the actual assembly situation of the RPM.

By analyzing the geometric relation of joints in PUU limbs, the equation of input variable $d_{4}$ can be obtained as follows:
$d_{4}= \pm \arccos \left(\frac{z+a c \theta s \beta}{\sqrt{(x+a s \theta s \beta)^{2}+(z+a c \theta s \beta)^{2}}}\right)$,
where the " $\pm$ " notation is the same as the " $\pm$ " notation of the output parameter $x$.

Equations (6) and (7) present the inverse position solutions for the proposed RPM. However, the four output parameters $(x, z, \theta, \varphi)$ are not totally independent since the 2PRU-PUU RPM has 3 DOFs in 1R2T and 2R1T operation modes. When it is in 1R2T operation mode, $\varphi$ is always zero. In the 2R1T operation mode, coordinates $x$ of the point $o^{\prime}$ will be changed with rotation around axis $y$, satisfying the equation
$x=z \tan \theta$.


Figure 2. Kinematics model of the 2PRU-PUU PM.

## 4 Dynamic analysis

In this paper, the kinematic equation of the mechanism is established from the viewpoint of energy using the Lagrange method, and the kinetic energy, potential energy, and generalized force of the system are analyzed; thus the dynamic equation of the RPM is derived, and the driving force of the mechanism is optimized.

### 4.1 Velocity analysis

The driving velocity vectors obtained by deriving Eqs. (6) and (7) with respect to the time derivative can be expressed as
$\dot{\boldsymbol{d}}=\mathbf{J}_{i} \dot{\boldsymbol{q}}$,
where $\dot{\boldsymbol{q}}=\left[\begin{array}{lll}\theta & x & z\end{array}\right]^{\mathrm{T}}$ in the 1R2T operation mode, and $\dot{\boldsymbol{q}}=\left[\begin{array}{lll}\theta & \varphi & z\end{array}\right]^{\mathrm{T}}$ in the 2R1T operation mode. $\mathbf{J}_{1}$ is the velocity Jacobian matrix of the RPM under the 1R2T operation mode, and $\mathbf{J}_{2}$ is the velocity Jacobian matrix of the RPM under the 2 R1T operation mode.
$\mathbf{J}_{1}=\left[\begin{array}{ccc}\frac{\Delta_{1}+a d_{1} s \theta+\Delta_{2}}{-\Delta_{3}} & -1 & \frac{-z-a s \theta}{\Delta_{3}} \\ \frac{\Delta_{1}-a d_{2} s \theta+\Delta_{2}}{-\Delta_{4}} & 1 & \frac{z-a s \theta}{\Delta_{4}} \\ 0 & \frac{x}{a-q_{3}} & \frac{z}{a-d_{3}} \\ 0 & \frac{x \cos \left(d_{4}\right)}{\sin \left(d_{4}\right) \Delta_{5}^{2}} & \frac{z \cos \left(d_{4}\right)}{\sin \left(d_{4}\right)\left(\Delta_{5}^{2}-\Delta_{5}\right)}\end{array}\right]$,
$\mathbf{J}_{2}=\left[\begin{array}{ccc}\frac{\Delta_{1} a d_{1} s \theta+\Delta_{2}}{-\Delta_{3}} & 0 & \frac{-z-a s \theta}{\Delta_{3}} \\ \frac{\Delta_{1}-a d_{2} s \theta+\Delta_{2}}{-\Delta_{4}} & 0 & \frac{z-a s \theta}{\Delta_{4}} \\ \frac{a s \varphi(z s \theta+x c \theta)}{d_{3}-a c \varphi} & \frac{\Delta_{6}}{a c \varphi-d_{3}} & \frac{z+a c \theta s \varphi}{a c \varphi-d_{3}} \\ 1 & 0 & 0\end{array}\right]$,
where $\quad \Delta_{1}=a z c \theta, \quad \Delta_{2}=\operatorname{axs} \theta, \quad \Delta_{3}=d_{1}+x-a c \theta$, $\Delta_{4}=x-d_{2}+a c \theta, \quad \Delta_{5}=\sqrt{x^{2}+z^{2}}, \quad$ and $\quad \Delta_{6}=$ $a\left(s \varphi\left(d_{3}+a s \theta\right)+c \varphi(z c \theta+x s \theta)\right)$.

The linear and angular velocity vectors of the moving platform can be expressed by independent generalized coordinates as
$\boldsymbol{V}_{\mathrm{p}}=\left[\begin{array}{ll}\boldsymbol{\omega}_{\mathrm{p}}^{\mathrm{T}} & \boldsymbol{v}_{\mathrm{p}}^{\mathrm{T}}\end{array}\right]^{\mathrm{T}}=\mathbf{K}_{i} \dot{\boldsymbol{q}}=\left[\begin{array}{ll}\mathbf{J}_{\mathrm{P}, \mathrm{R}}^{\mathrm{T}} & \mathbf{J}_{\mathrm{P}, \mathrm{T}}^{\mathrm{T}}\end{array}\right]^{\mathrm{T}} \dot{\boldsymbol{q}}$,
where $\mathbf{K}_{i}(i=1,2,3)$ represents the mapping matrix between the moving platform velocity and the independent generalized coordinate velocity under 1R2T operation mode, switch configuration, and 2R1T operation mode, respectively; they can be expressed as
$\mathbf{K}_{1}=\left[\begin{array}{llllll}0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1\end{array}\right]^{\mathrm{T}}$,
$\mathbf{K}_{2}=\left[\begin{array}{cccccc}0 & 1 & 0 & z \sec ^{2} \theta & 0 & 0 \\ c \theta & 0 & -s \theta & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \tan \theta & 0 & 1\end{array}\right]^{\mathrm{T}}$,
$\mathbf{K}_{3}=\left[\begin{array}{cccccc}0 & 1 & 0 & z \sec ^{2} \theta & 0 & 0 \\ c \theta & 0 & -s \theta & 0 & 0 & 0 \\ 0 & 0 & 0 & \tan \theta & 0 & 1\end{array}\right]^{\mathrm{T}}$.
The velocity $\mathbf{v}_{B_{i}}$ of universal joint points $B_{i}$ in a fixed coordinate frame can be expressed by the velocity of the moving platform as
$\boldsymbol{v}_{B_{i}}=\boldsymbol{v}_{\mathrm{p}}+\omega_{\mathrm{p}} \times \boldsymbol{r}_{B_{i}}(i=1,2,3)$.
The velocity vector $\boldsymbol{v}_{B_{i}}$ can also be expressed by the velocity of link as
$\boldsymbol{v}_{B_{i}}=\boldsymbol{v}_{p i}+\omega_{\mathrm{li}} \times \boldsymbol{n}_{i} L(i=1,2,3)$,
where $\boldsymbol{n}_{i}=\frac{\boldsymbol{r}_{B_{i}}-\boldsymbol{r}_{P_{i}}}{L}$ is the unit direction vector of the link $P_{i} B_{i}$.

By substituting Eq. (16) into Eq. (17) and multiplying the vector $\boldsymbol{n}_{i}$, the angular velocity of the link $P_{i} B_{i}$ can be expressed as
$\omega_{\mathrm{li}}=\frac{\boldsymbol{n}_{i} \times\left(\boldsymbol{v}_{\mathrm{p}}+\omega_{\mathrm{p}} \cdot \boldsymbol{r}_{B_{i}}-\boldsymbol{v}_{p i}\right)}{L}(i=1,2,3)$.
After obtaining the angular velocity of the link $P_{i} B_{i}$, the linear velocity at the center of mass of the link is
$\boldsymbol{v}_{\mathrm{li}}=\boldsymbol{v}_{p i}+\omega_{\mathrm{li}} \times \boldsymbol{n}_{i} \frac{L}{2}(i=1,2,3)$.

### 4.2 The kinetic energy of system

The kinetic energy of the RPM includes the kinetic energy of the moving platform, link, and driving pair. The kinetic energy of the moving platform can be expressed as
$T_{\mathrm{p}}=\frac{1}{2} m_{\mathrm{p}} \boldsymbol{v}_{\mathrm{p}}^{\mathrm{T}} \boldsymbol{v}_{\mathrm{p}}+\frac{1}{2} \omega_{\mathrm{p}}^{\mathrm{T}} \mathbf{I}_{\mathrm{p}} \boldsymbol{\omega}_{\mathrm{p}}=\frac{1}{2} \dot{\boldsymbol{q}}^{\mathrm{T}} \mathbf{M}_{\mathrm{p}} \dot{\boldsymbol{q}}$,
where $\mathbf{I}_{\mathrm{p}}$ represents the moment of inertia matrix at the center of mass of the moving platform.

The kinetic energy of three links of the RPM is
$T_{1}=\sum_{i=1}^{3}\left(\frac{1}{2} m_{\mathrm{li}} \boldsymbol{v}_{\mathrm{li}}^{\mathrm{T}} \boldsymbol{v}_{\mathrm{li}}+\frac{1}{2} \boldsymbol{\omega}_{\mathrm{li}}^{\mathrm{T}} \mathbf{I}_{\mathrm{li}} \boldsymbol{\omega}_{\mathrm{li}}\right)=\frac{1}{2} \dot{\boldsymbol{q}}^{\mathrm{T}} \mathbf{M}_{1} \dot{\boldsymbol{q}}$,
where $\mathbf{I}_{\mathrm{li}}$ represents the moment of the inertia matrix at the center of mass of the link.

Three sliders of 2PRU-PUU redundant actuated PMs only have translation motion along the coordinate axis, so there is no rotational kinetic energy generated by three sliders. Correspondingly, the fourth actuator has no linear kinetic energy but rotational kinetic energy.
$T_{\mathrm{d}}=\sum_{i=1}^{3} \frac{1}{2} m_{\mathrm{di}} \boldsymbol{v}_{\mathrm{di}}^{\mathrm{T}} \boldsymbol{v}_{\mathrm{di}}+\frac{1}{2} \boldsymbol{\omega}_{\mathrm{d} 4}^{\mathrm{T}} I_{\mathrm{d} 4} \boldsymbol{\omega}_{\mathrm{d} 4}=\frac{1}{2} \dot{\boldsymbol{q}}^{\mathrm{T}} \mathbf{M}_{\mathrm{d}} \dot{\boldsymbol{q}}$,
where $\boldsymbol{v}_{\mathrm{di}}$ denotes the linear velocity vectors of the driving slider, and $\omega_{\mathrm{d} 4}$ denotes the angular velocity vectors of the fourth actuator.
According to Eqs. (20), (21), and (22), the total kinetic energy of the system can be obtained as follows:
$T=T_{\mathrm{p}}+T_{1}+T_{\mathrm{d}}$.

### 4.3 Potential energy of system

The potential energy is related to the choice of coordinate frame. If the origin point $o$ is the zero potential energy point and the elasticity and friction of the components are ignored, the potential energy of the moving platform, link, and driving pair can be expressed as
$U_{\mathrm{p}}=m_{\mathrm{p}} z_{\mathrm{p}}^{\mathrm{T}} \mathbf{g}$,
$U_{1}=\sum_{i=1}^{3} m_{\mathrm{li}} z_{\mathrm{li}}^{\mathrm{T}} \mathbf{g}$,
$U_{\mathrm{d}}=\sum_{i=1}^{4} m_{\mathrm{di}} z_{\mathrm{d} \mathbf{i}}^{\mathrm{T}} \mathbf{g}$,
where $\mathbf{g}$ denotes the acceleration of gravity, and $z_{\mathrm{p}}, z_{\mathrm{li}}$, and $z_{\mathrm{di}}$ are the position vectors of the centroid of the moving platform, the links, and the driving pairs in the fixed coordinate frame.

The total potential energy of the system can be obtained from Eqs. (24), (25), and (26).
$U=U_{\mathrm{p}}+U_{\mathrm{l}}+U_{\mathrm{d}}$

### 4.4 Dynamic modeling

The dynamic models of the manipulator are now derived via the Lagrangian formulation as
$\frac{\mathrm{d}}{\mathrm{d} t}\left(\frac{\partial L}{\partial \dot{\boldsymbol{q}}}\right)-\frac{\partial L}{\partial \boldsymbol{q}}=\boldsymbol{\tau}$,
where $L=T-U$ is the Lagrangian function; $T$ and $U$ are the kinetic energy and potential energy functions of the system, respectively; and $\boldsymbol{\tau}$ is the generalized force.

The kinetic energy and potential energy of the system have been obtained in Sect. 4.2 and 4.3. By substituting Eqs. (23) and (27) into Eq. (28), the dynamic equation of 2PRU-PUU RPM can be obtained as follows
$\mathbf{M}(\boldsymbol{q}) \ddot{\boldsymbol{q}}+\boldsymbol{C}(\boldsymbol{q}, \dot{\boldsymbol{q}}) \dot{\boldsymbol{q}}+\boldsymbol{G}(\boldsymbol{q})=\boldsymbol{\tau}$,
where $\boldsymbol{\tau}=\left[\tau_{1}, \tau_{2}, \tau_{3}, \tau_{4}\right]^{\mathrm{T}}, \boldsymbol{G}(\boldsymbol{q})=\partial U / \partial \boldsymbol{q}$,
$\mathbf{M}(\boldsymbol{q})=\mathbf{M}_{\mathrm{p}}+\mathbf{M}_{\mathbf{l}}+\mathbf{M}_{\mathrm{d}}$,
$\boldsymbol{C}(\boldsymbol{q}, \dot{\boldsymbol{q}})=\left(\sum_{i=1}^{3} \dot{\boldsymbol{q}}_{i} \frac{\partial \mathbf{M}}{\partial \boldsymbol{q}_{i}}\right)-\frac{1}{2}\left[\dot{\boldsymbol{q}}^{\mathrm{T}} \frac{\partial \mathbf{M}}{\partial \boldsymbol{q}_{1}}, \dot{\boldsymbol{q}}^{\mathrm{T}} \frac{\partial \mathbf{M}}{\partial \boldsymbol{q}_{2}}, \dot{\boldsymbol{q}}^{\mathrm{T}} \frac{\partial \mathbf{M}}{\partial \boldsymbol{q}_{3}}\right]^{\mathrm{T}}$.
According to the principle of virtual work (Tsai, 1998), the relation between the generalized force and axial driving force can be expressed as
$\boldsymbol{\tau}=\mathbf{J}^{\mathrm{T}} \boldsymbol{F}$,
where $\boldsymbol{F}=\left[f_{1}, f_{2}, f_{3}, f_{4}\right]^{\mathrm{T}}$ is the driving force, and $\mathbf{J}^{\mathrm{T}} \in$ $\mathbf{R}^{3 \times 4}$ is the force Jacobian matrix of the system.

As we can see, there are three independent equations with four unknown driving forces in Eq. (29). However, $\mathbf{J}$ is a non-full-rank matrix, and the unique solution of the driving force cannot be solved directly. To obtain a unique solution, an optimization technique has to be applied. For different demands, different optimizing objectives are needed. In this paper, the optimization problem of the driving force of the weighted least-squares method can be expressed as follows: the optimization variable is the driving moment $\boldsymbol{F}$, and under the constraint condition $\boldsymbol{\tau}=\mathbf{J}^{\mathrm{T}} \boldsymbol{F}$, the driving force $\boldsymbol{F}$ of the objective function $Z=\boldsymbol{F}^{\mathrm{T}} \mathbf{W} \boldsymbol{F}$ is minimized, where $\mathbf{W}$ is a diagonal weighted matrix, which represents the weight of the driving force of each limb (Fontes and Da Silva, 2016; Park et al., 2003).

For the above optimization problems, the Lagrange multiplier method can be used, and the Lagrange multiplier $\lambda$ is introduced to construct a new function $Z^{\prime}$ :
$Z^{\prime}=\boldsymbol{F}^{\mathrm{T}} \mathbf{W} \boldsymbol{F}+\lambda^{\mathrm{T}}\left(\boldsymbol{\tau}-\mathbf{J}^{\mathrm{T}} \boldsymbol{F}\right)$.
In order to get the extremum of $Z^{\prime}$, the following conditions must be satisfied

$$
\left\{\begin{array}{l}
\frac{\partial Z^{\prime}}{\partial \boldsymbol{F}}=2 \boldsymbol{F}^{\mathrm{T}} \mathbf{W}-\lambda^{\mathrm{T}} \mathbf{J}^{\mathrm{T}}=0  \tag{34}\\
\frac{\partial Z^{\prime}}{\partial \lambda}=\tau-\mathbf{J}^{\mathrm{T}} \boldsymbol{F}=0
\end{array}\right.
$$

From Eq. (34), the driving force $\boldsymbol{F}$ can be expressed as

$$
\begin{equation*}
\boldsymbol{F}=\left(\mathbf{W}^{-1}\right)^{\mathrm{T}} \mathbf{J}\left(\mathbf{J}^{\mathrm{T}}\left(\mathbf{W}^{-1}\right)^{\mathrm{T}} \mathbf{J}\right)^{-1} \boldsymbol{\tau} \tag{35}
\end{equation*}
$$

Taking the weighting matrix $\mathbf{W}$ as the unit matrix,
$\boldsymbol{F}=\mathbf{J}\left(\mathbf{J}^{\mathrm{T}} \mathbf{J}\right)^{-1} \boldsymbol{\tau}$.

### 4.5 Verification with physical simulation

It is known that $a=0.09 \mathrm{~m}, \quad m_{\mathrm{p}}=0.254 \mathrm{~kg}, \quad m_{11}=$ $m_{12}=0.238 \mathrm{~kg}, \quad m_{13}=0.393 \mathrm{~kg}, \quad m_{\mathrm{d} 1}=m_{\mathrm{d} 2}=0.053 \mathrm{~kg}$, $m_{\mathrm{d} 3}=0.041 \mathrm{~kg}$, and $L=0.33 \mathrm{~m}$. The correctness of the above dynamic models is verified by ADAMS dynamic simulation software. The trajectory equation of 1R2T operation mode is

$$
\left\{\begin{array}{l}
\theta=\left(10 \pi-10 \pi\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right)\right) / 180  \tag{37}\\
x=-30+30\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right) \\
z=250-20\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right)
\end{array}\right.
$$

The trajectory equation of 2 R 1 T operation mode is

$$
\left\{\begin{array}{l}
\theta=10 \pi\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right) / 180  \tag{38}\\
\varphi=10 \pi\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right) / 180 \\
z=230+20\left(6 t^{5} / t_{\mathrm{d}}^{5}-15 t^{4} / t_{\mathrm{d}}^{4}+10 t^{3} / t_{\mathrm{d}}^{3}\right)
\end{array}\right.
$$

where $t_{\mathrm{d}}=5$ denotes the duration of the motion.
To address the issue that the inverse dynamics solution in the simulation process of redundant actuated PMs is not unique, the displacement of the actuating joints from limb 1 to limb 3 and the theoretical value of the driving 4 are input into ADAMS software, and the unique driving force combination can be obtained.

In Fig. 3, the periods of $1-5$ and 5-10 s can be compared relating to the theoretical results and the simulation results of the three limbs' driving force of 1R2T and 2R1T operation mode, respectively. It is evident that the theoretical results and simulation results produce the same computed force outputs, which indicates the correctness of the theoretical analysis, and the dynamic models can be used for the subsequent dynamic performance evaluation.

## 5 Dynamic performance analysis

The dynamic manipulability ellipsoid (DME) index, which is frequently used to investigate the dynamic performance of serial or parallel manipulators, can be used to evaluate the ability of the mechanism to change the position or orientation of the end effector under the constraint of the driving force (Chai et al., 2020). The approximate mapping relation between the generalized acceleration vector and the driving force vector can be obtained by ignoring the velocity factor and the gravity factor. Equation (36) can be rewritten as
$\boldsymbol{F}=\mathbf{J}^{-\mathrm{T}} \mathbf{H}^{\mathrm{T}} \boldsymbol{\Phi} \mathbf{H} \mathbf{J}_{\mathrm{P}}^{+} \boldsymbol{a}_{\mathrm{P}}$,
where $\mathbf{H}$ represents the pose matrix of each component of the mechanism, $\boldsymbol{\Phi}$ represents the inertia tensor and mass matrix of each component in the mechanism, $\mathbf{J}_{\mathrm{P}}=\left[\mathbf{J}_{\mathrm{P}, \mathrm{R}}^{\mathrm{T}}, \mathbf{J}_{\mathrm{P}, \mathrm{T}}^{\mathrm{T}}\right]^{\mathrm{T}}$ is the Jacobian matrix between the moving platform and generalized coordinates, and $\boldsymbol{a}_{\mathrm{p}}$ is the acceleration of the moving platform.


Figure 3. Verification with physical simulation.

According to Eq. (39), an ellipsoid can be obtained to describe the acceleration boundary of the end effector in various directions under the constraint of driving force as follows:

$$
\begin{equation*}
\|\tilde{\boldsymbol{F}}\| \leq 1 \Rightarrow \tilde{\boldsymbol{a}}^{\mathrm{T}} \tilde{\mathbf{J}}^{+\mathrm{T}} \tilde{\mathbf{M}}^{\mathrm{T}} \tilde{\mathbf{M}} \tilde{\mathbf{J}}^{+} \tilde{\boldsymbol{a}} \leq 1, \tag{40}
\end{equation*}
$$

where $\tilde{\boldsymbol{F}}$ represents the generalized driving forces, $\tilde{\mathbf{M}}=$ $\mathbf{J}_{i}^{-\mathrm{T}} \mathbf{H}^{\mathrm{T}} \boldsymbol{\Phi} \mathbf{H} \mathbf{J}_{i}^{-1}$ is the inertia matrix of the mechanism, and $\tilde{\boldsymbol{a}}$ represents the acceleration of the end effector. $\tilde{\mathbf{J}}=\mathbf{J}_{\mathrm{P}} \mathbf{J}_{i}^{-1}$ is the Jacobian matrix relating the actuated joints of mechanism to its end effector.

By replacing the corresponding Jacobian matrices with separate Jacobian matrices (Chen et al., 2017), the DME index for rotational and translational of the mechanism can be derived from Eq. (40), which is expressed as

$$
\|\tilde{\boldsymbol{F}}\| \leq 1 \Rightarrow\left\{\begin{array}{l}
\tilde{\boldsymbol{a}}_{\mathrm{R}}^{\mathrm{T}} \tilde{\mathbf{J}}_{\mathrm{R}}^{+\mathrm{T}} \tilde{\mathbf{M}}^{\mathrm{T}} \tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{R}}^{+} \tilde{\boldsymbol{a}}_{\mathrm{R}} \leq 1  \tag{41}\\
\tilde{\boldsymbol{a}}_{\mathrm{T}}^{\mathrm{J}} \tilde{\mathbf{J}}_{\mathrm{T}}^{+T} \tilde{\mathbf{M}}^{\mathrm{T}} \tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{T}}^{+} \tilde{\boldsymbol{a}}_{T} \leq 1,
\end{array}\right.
$$

where $\tilde{\mathbf{J}}_{\mathrm{R}}=\mathbf{J}_{\mathrm{P}, \mathrm{R}} \mathbf{J}_{i}^{-1}$ and $\tilde{\mathbf{J}}_{\mathrm{T}}=\mathbf{J}_{\mathrm{P}, \mathrm{T}} \mathbf{J}_{i}^{-1}$ correspond to the Jacobian matrix of rotational and translational, respectively.

The isotropic property of rotational dynamic manipulability of the 2PRU-PUU RPM is evaluated by the $\tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{R}}^{+}$condition number, which is expressed as
$\omega_{\mathrm{R}}=\frac{\sigma_{\mathrm{R} 1}}{\sigma_{\mathrm{R} 2}}$,
where $\boldsymbol{\sigma}_{\mathrm{R} 1}$ and $\boldsymbol{\sigma}_{\mathrm{R} 2}$ are nonzero singular values of $\tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{R}}^{+}$, and $\sigma_{\mathrm{R} 1} \geq \sigma_{\mathrm{R} 2}$.

In the same way, the isotropic property of translational dynamic manipulability is evaluated by the $\tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{T}}^{+}$condition number, which is expressed as
$\omega_{\mathrm{T}}=\frac{\boldsymbol{\sigma}_{\mathrm{T} 1}}{\boldsymbol{\sigma}_{\mathrm{T} 2}}$,
where $\boldsymbol{\sigma}_{\mathrm{T} 1}$ and $\boldsymbol{\sigma}_{\mathrm{T} 2}$ are nonzero singular values $\tilde{\mathbf{M}} \tilde{\mathbf{J}}_{\mathrm{T}}^{+}$, and $\sigma_{\mathrm{T} 1} \geq \boldsymbol{\sigma}_{\mathrm{T} 2}$.

According to Eqs. (42) and (43), the performance distribution of dynamic manipulability of rotation and translation
of 2PRU-PUU RPM in two operation modes at $z=230 \mathrm{~mm}$ can be obtained, as shown in Figs. 4 and 5. In Fig. 4, the rotational and translational dynamic manipulability of 1R2T operation mode is symmetrical with respect to $\theta=6.84^{\circ}$ because the 2PRU-PUU RPM is symmetrical with respect to $\theta=6.84^{\circ}$ in the 1 R 2 T operation mode, and the mechanism is symmetrical with respect to $\theta=0^{\circ}$ when moving to the 2R1T operation mode. As can be seen from Fig. 5, the rotational and translational dynamic manipulability of 2R1T operation mode is symmetrical with respect to $\theta=0^{\circ}$, which conforms to the structural characteristics of the 2PRU-PUU RPM. At $\theta \in\left(-10^{\circ}, 10^{\circ}\right)$ the isotropic property becomes better because the effect of the differential motion between limb 1 and limb 2 decreases with a more oblique configuration of the end effector.

## 6 Conclusions

This paper presents kinematic, dynamic, and performance analysis of a novel 2PRU-PUU RPM. From the switch configuration, the mechanism can evolve into 1R2T operation mode and 2R1T operation mode. The dynamic models of the 2PRU-PUU RPM are established based on the Lagrange method. To solve the problem of the non-uniqueness solution of redundant actuator PMs in dynamic analysis, the weighted least-squares norm is introduced. The results of the modeling method have been validated by means of numerical simulations of commercial ADAMS software. The DME index is then adopted to evaluate the dynamic manipulability performance of the 2PRU-PUU RPM. And the dynamic performance of rotation and translation of the mechanism in the workspace is derived. The overall performance of the two operation modes of the RPM is better at the operating height of $z=0.23 \mathrm{~m}$ and $\theta \in\left(-10^{\circ}, 10^{\circ}\right)$. It provides a basis for subsequent performance optimization and prototype motion control. In future work, we will investigate the multi-objective optimization design of the 2PRU-PUU RPM.


Figure 4. Dynamic performance distribution of 1R2T operating mode. (a) Rotational dynamic manipulability index. (b) Translational dynamic manipulability index.


Figure 5. Dynamic performance distribution of 2R1T operating mode. (a) Rotational dynamic manipulability index. (b) Translational dynamic manipulability index.

Data availability. The data generated during this study are available from the corresponding author on reasonable request.

Author contributions. Conceptualization: TS and CY. Methodology: WY and FH. Software: TS and CY. Validation: TS and FH. Formal analysis: WY and TS. Investigation: TS and CY. Resources: FH. Writing (original draft preparation): FH and CY. Writing (review and editing): TS and CY. Visualization: WY and TS. Supervision: WY. Project administration: FH. Funding acquisition: WY and FH.

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